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## Wind-Forced Reversing Jets in the Western Equatorial Pacific\*

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### ABSTRACT

Upper-ocean zonal currents in the western equatorial Pacific are remarkably variable, changing direction both with time and depth. As a part of the Tropical Ocean and Global Atmosphere Coupled Ocean–Atmosphere Response Experiment, an enhanced monitoring array of moorings measured the upper-ocean velocity, temperature, salinity, and, surface meteorological conditions in the western equatorial Pacific for two years (March 1992–April 1994). Data from this array are used to evaluate the zonal momentum balance. Although nonlinear terms (zonal, meridional, and vertical advection) were at times large, reversing jets were primarily due to an interplay between wind forcing and compensating pressure gradients. In the weakly stratified surface layer, the flow is to a large extent directly forced by local winds. Eastward acceleration associated with westerly wind bursts and westward accelerations associated with easterly trades lead to frequent reversals in the surface-layer flow. However, pressure gradients set up by the wind bursts partially compensate the local wind forcing in the surface layer. Below the surface layer, these pressure gradients tend to accelerate the upper-thermocline flow in a direction opposing the local winds. Consequently, during westerly wind bursts, a reversing jet structure can develop, with a surface eastward current overlying a westward intermediate layer flow, overlaying the eastward Equatorial Undercurrent.

### 1. Introduction

Wind forcing has a particularly dramatic ocean response on the equator. With no Coriolis turning, local accelerations can be large, leading to strong currents. However, even for sustained wind forcing, equatorial currents do not accelerate indefinitely. Because of the relatively fast equatorially trapped waves, the equatorial ocean can adjust rapidly to wind forcing by setting up an opposing pressure gradient. The resulting pressure gradient can support a subsurface counterflow, as for example in the case of the eastward Equatorial Undercurrent (EUC) beneath the surface westward South Equatorial Current (SEC) for trade wind forcing. In this paper, we focus on the ocean response to wind forcing in the western equatorial Pacific where frequent reversals in the trade winds associated with westerly wind bursts

(WWBs) can lead to complex reversing current structures in the upper 300 m.

Reversing currents associated with westerly wind bursts were first observed by Hisard et al. (1970). During a cruise along 170°E in March 1967, trade winds were prevailing and the equatorial currents had a two-layer structure, with a westward 40 cm s<sup>-1</sup> SEC from the surface to 60 m, and an eastward EUC from 60 m to 300 m. One month later, the currents had a three-layer structure, with an eastward 20 cm s<sup>-1</sup> flow in the upper 60 m, westward 40 cm s<sup>-1</sup> flow between 60 m and 175 m, and eastward EUC flow between 175 m and 300 m. In the intervening month between cruises, a WWB had occurred.

Since the Hisard et al. (1970) study, there have been numerous reports of reversing jets in the western equatorial Pacific (e.g., McPhaden et al. 1988, 1990, 1992; Kuroda and McPhaden 1993; Delcroix et al. 1993; Smyth et al. 1996), and even as far east as 170°W during El Niños (Weisberg and Hayes 1995; Weisberg and Wang 1997). Due to the lack of data, dynamical analyses of the reversing jet have been incomplete. Using a combination of shipboard, island, and moored time series, McPhaden et al. (1988) showed that during the first several days of the WWB, an eastward “Yoshida jet” (Yoshida 1959) formed with surface acceleration rough-

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ly corresponding to the zonal wind stress forcing. A Kelvin wave generated by the WWB could subsequently be seen propagating eastward across the basin. Thus they hypothesized that the discrepancy between wind stress and surface acceleration after the first few days of the wind burst was due to the pressure gradient variability associated with the Kelvin wave passage. Using hydrographic data from zonal equatorial transects, McPhaden et al. (1990) later observed pressure gradient variability associated with WWBs.

Turbulent mixing is also expected to be important. In particular, McPhaden et al. (1988) apply the Stommel (1960) homogeneous EUC model to the surface layer to explain the development of vertical shear in terms of mixing. Microstructure measurements taken at 2°S, 156°E during a WWB (Smyth et al. 1996) show that turbulent kinetic energy increased during the WWB and, while it was highest in the surface mixed layer, patches of high turbulent kinetic energy were also observed on the upper and lower boundaries of the subsurface westward jet between the surface wind-driven eastward jet and EUC.

Scale analyses suggest that nonlinearities also may be important during strong westerly wind events due to zonal convergence, associated with the Yoshida jet impinging on the unforced region, and meridional Ekman convergence and downwelling; however, advective acceleration has been measured on synoptic timescales only in the central equatorial Pacific (Qiao and Weisberg 1997). There, advective acceleration at any individual depth was found to be comparable in magnitude with the local acceleration, although upon vertical integration across the EUC, the advective acceleration terms tended to cancel.

Unlike the steady trade winds found in the central and eastern tropical Pacific, trade winds in the western equatorial Pacific tend to be weak and punctuated by westerly wind bursts. Harrison and Vecchi (1997) define three westerly wind event composites for the equatorial Pacific, two of which are centered west of the date line (0°, 142.5°E; 0°, 170°E; 2.5°S, 170°W). The composites had peak wind speeds of 5.5–6.9 m s<sup>-1</sup>, fetch 3400–3800 km, duration 6–7 days, and were either stationary or propagate very slowly westward at 1.3 m s<sup>-1</sup>. Westerly wind bursts in the western equatorial Pacific often occur during the active convective phase of the 30–60 day intraseasonal Madden-Julian oscillation (MJO) (Madden and Julian 1994), which tend to occur more frequently during November through March and during El Niño years. Superclusters of convective clouds associated with MJO have been observed to have an eastward propagation of up to 10–15 m s<sup>-1</sup> (Nakazawa 1988).

Analytical and numerical models of the ocean response to wind forcing (e.g., Cane and Sarachik 1976; McCreary 1985; Tang and Weisberg 1984; McCreary and Lukas 1986; Weisberg and Tang 1987; Eriksen 1993; Giese and Harrison 1990; Weisberg and Tang

1990; Harrison and Craig 1993; Zhang and Rothstein 1998; Richardson et al. 1999) show that these subtle differences in the WWB characteristics can cause quite different ocean responses. In particular, wind bursts translating at speeds of 10 m s<sup>-1</sup> predominately excite inertia gravity waves (Erikson 1993), while eastward and westward slowly translating wind bursts tend to excite respectively near-resonant Kelvin or Rossby waves (Weisberg and Tang 1983; Tang and Weisberg 1984; McCreary and Lukas 1986). Recently, Richardson et al. (1999) simulated the formation of a subsurface westward reversing jet by forcing an ocean general circulation model with a stationary WWB. Analysis of the current structure evolution showed that the subsurface westward accelerations were primarily due to the propagation and separation of baroclinic Rossby and Kelvin waves generated at the edges, particularly the eastern edge, of the wind patch.

In this paper, data from the Tropical Ocean Global Atmosphere Coupled Ocean–Atmosphere Response Experiment (TOGA COARE) enhanced monitoring array of Acoustic Doppler Current Profilers (ADCPs) and Tropical Atmosphere Ocean (TAO) thermistor moorings with SEACAT conductivity (salinity) sensors and surface meteorological packages are used to analyze the variability and dynamics of the wind-forced upper-ocean zonal current structures in the western equatorial Pacific. With this dataset, nearly all terms in the zonal momentum balance can be evaluated for a period of up to two years. Although the basic westward South Equatorial Current overlying the eastward Equatorial Undercurrent is often seen in the data, the vertical structure is highly variable and also exhibits periods with all eastward flow, as well as periods with reversing jets with complex vertical structure. These stacked “ribbons” of flow, sometimes only 50 m thick, are often coherent over several weeks and over the full zonal extent of our array (1800 km). Based on observations by Kuroda and McPhaden (1993), these reversing jets can at times be coherent even over ranges of up to a few thousand kilometers. The outline of the paper is as follows: In the next section, data used in the analysis are discussed. In section 3, the methodology is presented. Our main analysis tool is the zonal momentum balance. Results are presented in section 4 and summarized and discussed in section 5.

## 2. Data

The TAO array consists of nearly 70 Autonomous Temperature Line Acquisition System (ATLAS) moorings in the tropical Pacific and five current meter moorings along the equator. All moorings measure daily averaged wind speed and direction, air temperature, and relative humidity, and surface and subsurface temperature from 1 m to 500 m with at least a 1-h sample rate (McPhaden et al. 1998). As a part of TOGA COARE, the TAO array was enhanced with additional moorings

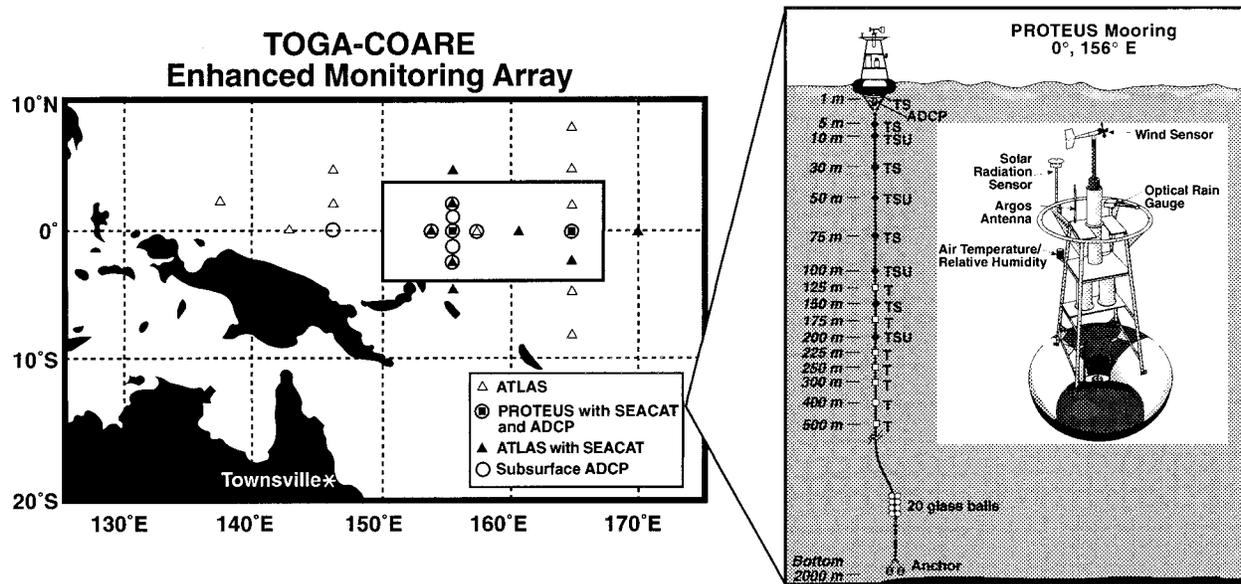


FIG. 1. The COARE Enhanced Monitoring Array (EMA). The array included moored ADCPs, Tropical Atmosphere Ocean (TAO) ATLAS moorings, some of which were enhanced with SEACAT temperature and conductivity sensors, and TAO PROTEUS moorings enhanced with SEACAT sensors. The  $0^{\circ}$ ,  $156^{\circ}\text{E}$  PROTEUS mooring is shown in the inset.

and instrumentation to monitor upper-ocean currents, salinity, and temperature (with higher spatial and temporal resolution) along the equator and the  $156^{\circ}\text{E}$  meridian for approximately two years (March 1992–April 1994). The enhanced monitoring array (EMA) is shown in Fig. 1, and schematic diagrams of the data availability are shown in Fig. 2. At the  $0^{\circ}$ ,  $156^{\circ}\text{E}$  central site, nearly all terms in the zonal momentum balance could be evaluated. These data and the surface meteorological data used in the analysis are described below.

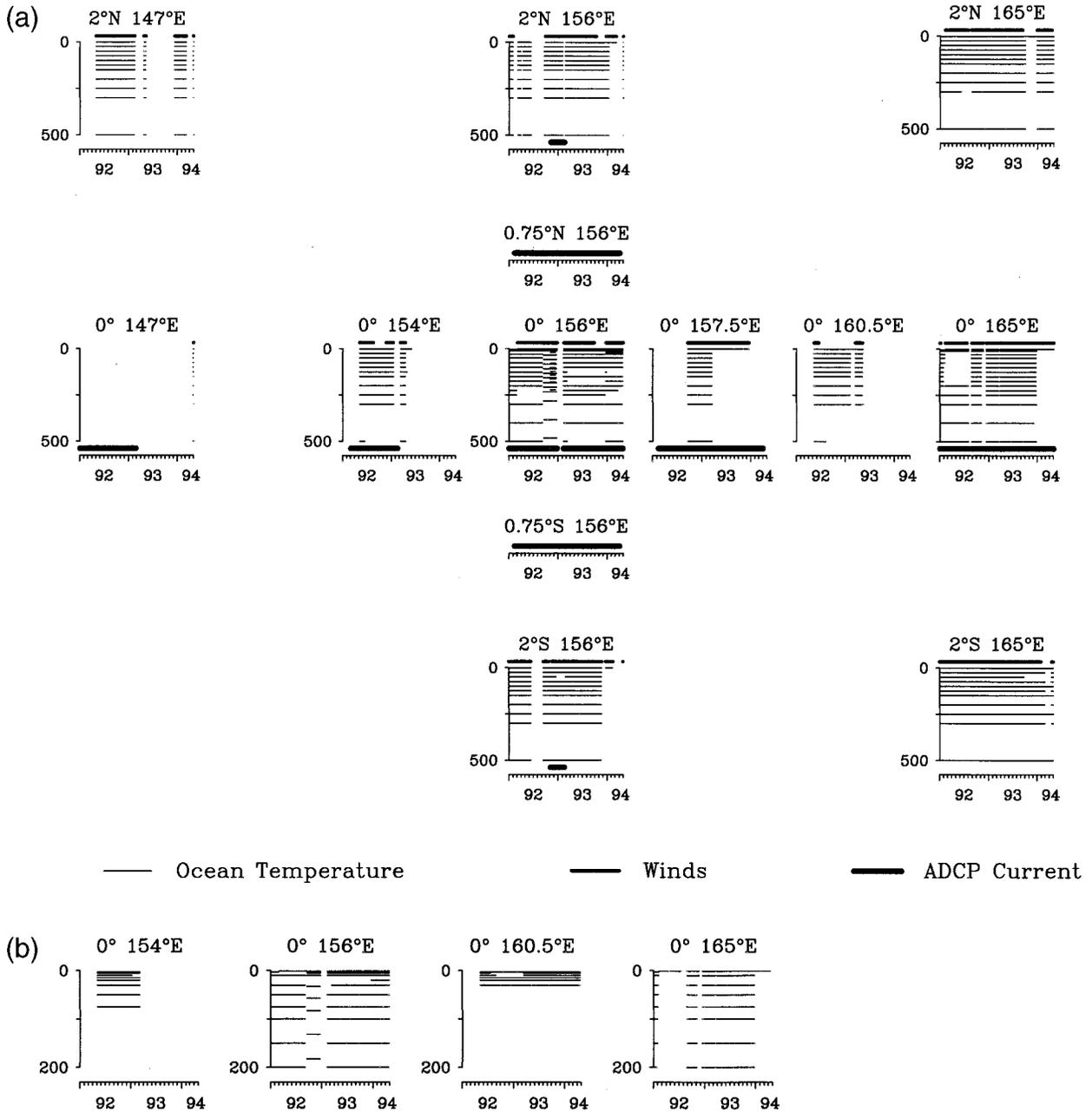
#### a. Currents

Upper-ocean currents in the EMA region were measured at five sites along the equator and four sites off the equator along the  $156^{\circ}\text{E}$  meridian (Figs. 1, 2, Table 1) by an array of 153.6-kHz ADCPs. Two sites,  $0^{\circ}$ ,  $165^{\circ}\text{E}$  and  $0^{\circ}$ ,  $156^{\circ}\text{E}$  (from August 1991 to March 1993), had Profile Telemetry of Upper Ocean Currents (PROTEUS) moorings with downward-looking ADCPs mounted on the buoy's bridle and several mechanical current meters attached to the mooring line. All other ADCPs were upward-looking subsurface-moored ADCPs. Figure 3 shows the zonal ADCP currents along the equator, and Fig. 4a shows the meridional ADCP currents at the central  $0^{\circ}$ ,  $156^{\circ}\text{E}$  site. The ADCPs measured currents with at least hourly temporal resolution and 8–10-m vertical resolution. The depths of the top and bottom bins (Table 1) depend upon the geometry of the ADCP placement, but typically were 14–30 m and 230–280 m, respectively. ADCPs assume that the sound reflectors are small particles and phytoplankton

moving passively with the currents. Fish swarming near the surface buoys are actively swimming reflectors and thus can bias downward-looking ADCP data. However, analysis of the ADCP echo intensities and comparisons with the mechanical current meters indicate that fish biases were minimal at COARE sites (Plimpton et al. 1997).

For the purposes of our analyses, currents were extrapolated to the surface by assuming constant flow above the top measurement level. To fill gaps in the zonal current data at the central site ( $0^{\circ}$ ,  $156^{\circ}\text{E}$ ), current data from the four nearby sites ( $0^{\circ}$ ,  $154^{\circ}\text{E}$ ;  $0^{\circ}$ ,  $157.5^{\circ}\text{E}$ ;  $0.75^{\circ}\text{N}$ ,  $156^{\circ}\text{E}$ ;  $0.75^{\circ}\text{S}$ ,  $156^{\circ}\text{E}$ ) were used to generate a frequency-dependent transfer function. The filled data are valid for timescales greater than approximately 5 days and thus the filled data were smoothed with a 5-day triangular filter (3-day cutoff, 7-day half amplitude) and subsampled once per day.

One objective of the ADCP array was the computation of vertical velocity based on estimations of the horizontal divergence at the central  $0^{\circ}$ ,  $156^{\circ}\text{E}$  site (Helber and Weisberg 1998). From conservation of mass, a horizontal divergence implies vertical velocity convergence. Thus vertical velocity (Fig. 4b) can be computed by vertically integrating the horizontal divergence estimate and assuming the boundary condition of no vertical flow at the surface (e.g., Wyrtki 1981; Halpern and Freitag 1987; Weingartner and Weisberg 1991; Helber and Weisberg 1998; Weisberg and Qiao 1999). Note that during the second year (February 1993 through March 1994), zonal divergence was estimated with a side finite difference ( $0^{\circ}$ ,  $157.5^{\circ}\text{E}$  minus  $0^{\circ}$ ,  $156^{\circ}\text{E}$ ). All other zonal



SEACAT Temperature and Salinity

FIG. 2. (a) Data availability of EMA wind, subsurface temperature, and ADCP currents. The standard TAO depths for subsurface temperature in this region are 1 m, 25 m, 50 m, 75 m, 100 m, 125 m, 150 m, 200 m, 250 m, 300 m, and 500 m. Some moorings had enhanced vertical resolution for temperature beyond these standard depths. (b) Data availability of EMA SEACAT temperature and salinity along equator.

and meridional divergence estimates used central finite differences.

b. Winds

Each TAO buoy had as standard instrumentation an RM Young wind sensor and vane mounted atop the buoy

tower at 4 m above the ocean surface (Fig. 1). Wind speed and direction were sampled at 2 Hz and vector averaged over six minutes once per hour. As in Cronin and McPhaden (1997), wind stress was computed using hourly surface meteorological data at 0°, 156°E (wind speed and direction at 4 m, air temperature and relative humidity at 3 m, shortwave radiation when available,

TABLE 1. ADCP placement. Upward and downward looking ADCPs are specified as u and d in column 3. Subsurface moored upward-looking ADCPs were anchored typically 5-km from the corresponding surface mooring site. 153.6-kHz ADCPs measure horizontal currents in 8–10-m bins. The center depth of the top and bottom bins are listed in columns 4 and 5. Institutions listed in column 6 were responsible for collecting and processing the ADCP data. The upward-looking ADCP at 0°, 156°E was loaned to PMEL by USF.

Site	Date	Orient	Top bin (m)	Bottom bin (m)	Institution
0°, 147°E	1 Jan 92–23 Feb 93	u	50	230	Tokai U.
0°, 154°E	1 Mar 92–16 Dec 93	u	20	240	Tokai U.
0°, 156°E	29 Aug 91–8 Mar 93	d	14	230–255	PMEL
0°, 156°E	12 Mar 93–27 Apr 94	u	20	260	PMEL
0°, 157.5°E	12 Feb 92–31 Mar 94	u	30	280	USF
0°, 165°E	8 Aug 91–16 Apr 94	d	14	220–240	PMEL
0.75°N, 156°E	12 Feb 92–3 Apr 94	u	30	280	USF
0.75°S, 156°E	12 Feb 92–2 Apr 94	u	30	280	USF

and SST at 1 m) and the COARE v2.5b bulk flux algorithm (Fairall et al. 1996a,b). The algorithm follows the standard Monin–Obukhov similarity approach for near-surface meteorological measurements, but includes models for the ocean’s cool skin and diurnal warm layer to estimate the skin temperature from measurements of the bulk SST at 1-m depth. The 5-day triangular filtered wind stress time series at 0°, 156°E is shown in Fig. 5.

### c. Subsurface temperature, salinity, and dynamic height

Dynamic height depends upon a vertical integral of temperature and salinity above some reference level. As shown in Fig. 2a, TAO moorings had daily averaged temperature data at depths ranging from 1 m to 500 m. As shown in Fig. 2b, enhancements included Seabird SEACAT temperature and conductivity sensors (from which salinity is computed) down to 200 m at 0°, 156°E and 0°, 165°E and down to only 70 m at 0°, 154°E and 30 m at 0°, 160.5°E (Cronin and McPhaden 1998; Freitag et al. 1999). To extrapolate the salinity profiles to the temperature measurement levels, a slowly varying, spatially dependent temperature–salinity ( $T$ – $S$ ) curve was computed by blending CTD and buoy temperature and salinity measurement pairs using an optimal interpolation scheme described in the appendix. Quality control and postprocessing of SEACAT data are described by Cronin and McPhaden (1998) and Freitag et al. (1999).

These daily averaged temperature and salinity data were then used to compute dynamic height relative to 500 db. At some sites the bottom thermistor(s) failed and the deepest thermistor was at 250 m or 300 m (Fig. 2). For these periods, the 500-db referenced dynamic height was computed based on a regression analysis, for example, 300 db versus 500 db referenced dynamic heights. The regressed dynamic heights typically had a root-mean-square error of roughly 0.6 dynamic centimeters at the surface.

In order to investigate the variability in the pressure gradient due to salinity, dynamic height was also computed using only temperature data [a climatological  $T$ – $S$  curve (Levitus and Boyer 1994; Levitus et al. 1994) was

used to estimate salinity from temperature]. The differences between the dynamic height estimates are due to variations in the  $T$ – $S$  curve and can be interpreted as the dynamic height due to anomalous variability in salinity. At 0°, 156°E the mean, rms, maximum, and minimum differences between the surface dynamic heights computed with and without observed salinity were 5, 6, 10, and  $-2$  dyn cm, respectively. These differences are considerable in relation to the 6 dyn cm standard deviation of the surface dynamic height at 0°, 156°E. However, since changes in the  $T$ – $S$  curve are often coherent over several degrees of longitude, salinity variations typically have a less dramatic effect on the zonal pressure gradient. The effects of salinity variations on pressure gradients and the momentum balance are analyzed further in section 3.

### 3. Methodology

In contrast to the SEC and EUC current structure found in the central equatorial Pacific, the western equatorial Pacific zonal currents show complex and variable vertical structures (Fig. 3) with current reversals both in depth and time. The depths of current reversals (i.e., zero current) roughly correspond to the depths of the 28°C and 20°C isotherms. Note that the 28°C isotherm is near the top of the thermocline and the location at which it outcrops defines the boundary of the western equatorial Pacific “warm pool.” Likewise, the 20°C isotherm is near the middle of the thermocline and its variability is often used as a proxy for thermocline depth variability (e.g., Rebert et al. 1985; Kessler et al. 1995). Thus we define three layers: a “surface layer” above the top of the thermocline from the surface to the depth of the 28°C isotherm (or 25 m, whichever is deeper), an “intermediate layer” in the upper thermocline between the base of the surface layer and the 20°C isotherm, and an “Equatorial Undercurrent layer” between the 20°C isotherm and the bottom ADCP bin (240 m). Because of the large variability within the surface layer, the record-length mean currents were near zero above 100 m (Fig. 3). In contrast, within the EUC layer, the currents were always eastward.

In order to understand the dynamics that give rise to

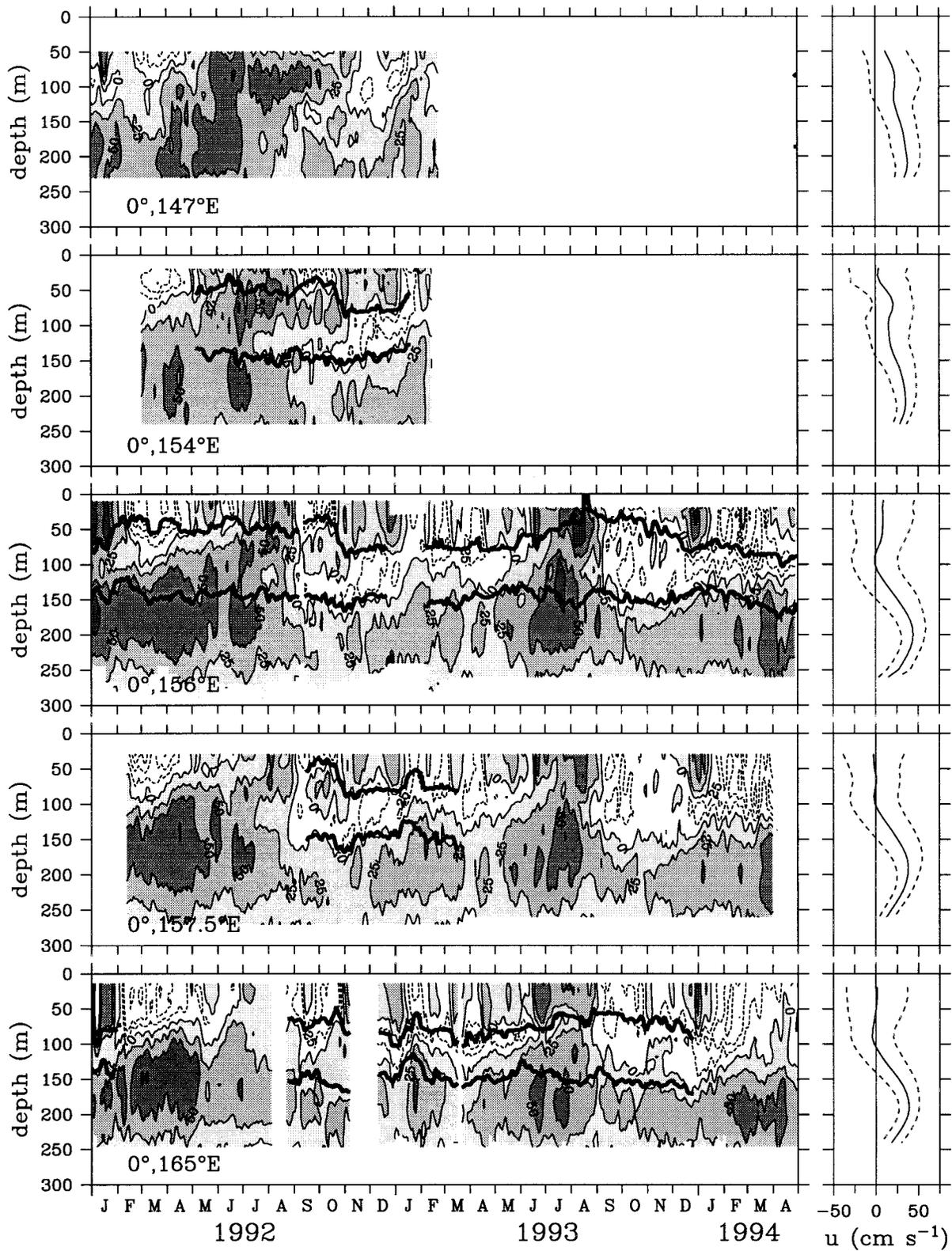


FIG. 3. Five-day triangular filtered zonal currents along the equator at 147°E, 154°E, 156°E, 157.5°E, and 165°E. The contour interval (CI) is 25 cm s<sup>-1</sup>. Eastward flow is shaded. The upper and lower thick lines are, respectively, 28°C and 20°C isotherms as measured by the corresponding TAO mooring. Right panels show the mean zonal currents and standard deviation envelope for each site along the equator.

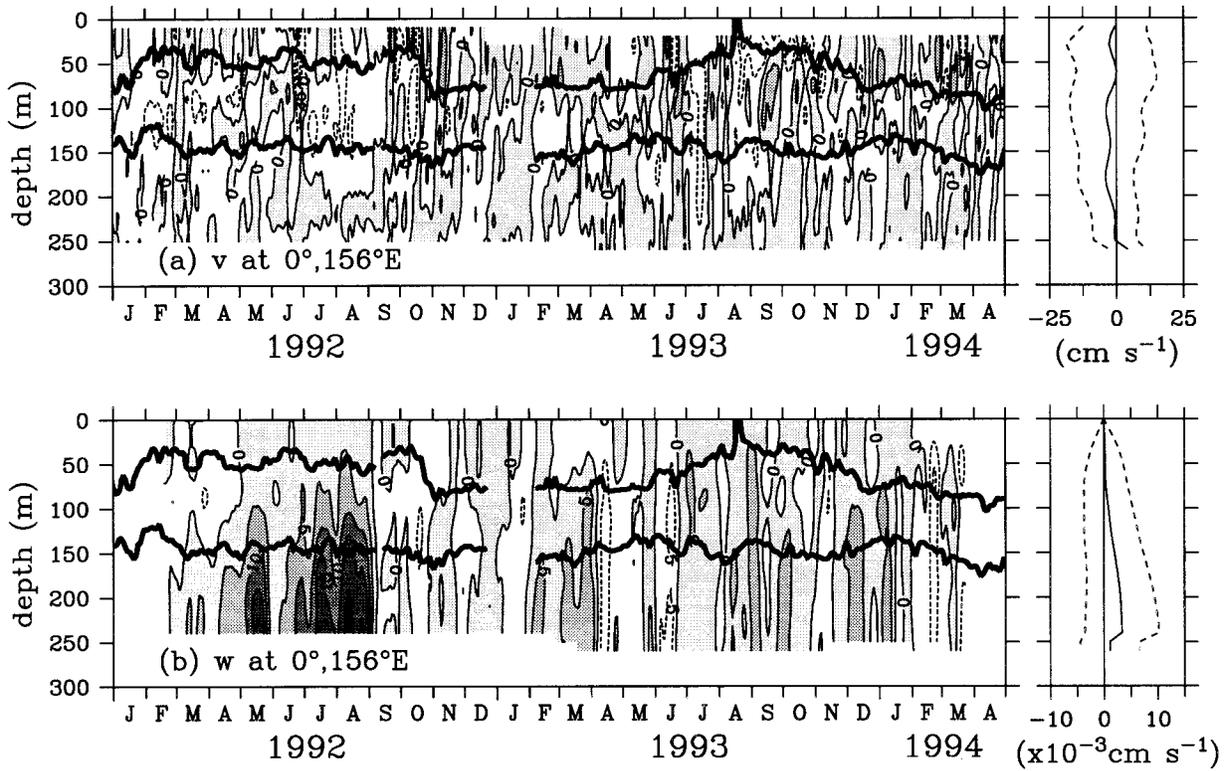


FIG. 4. (a) Five-day filtered meridional velocity as measured by the 0°, 156°E ADCP; CI is 25 cm s<sup>-1</sup>. Northward flow is shaded. Right panel shows the mean and standard deviation envelope of the 5-day filtered meridional flow. (b) Fifteen-day triangular filtered vertical velocity at 0°, 156°E as computed from mass budget in units 10<sup>-3</sup> cm s<sup>-1</sup>. The CI is 5 × 10<sup>-3</sup> cm s<sup>-1</sup>. Upwelling is shaded. Right panel shows the mean and the standard deviation envelope of the 5-day-filtered vertical velocity. The upper and lower thick lines in the left panels are, respectively, the 28°C and 20°C isotherms at 0°, 156°E.

these layered jet structures, the zonal momentum balance is evaluated at the 0°, 156°E central site both as a function of depth and within the three different layers. Forcing whose sign and magnitude change in successively deeper layers can cause complex vertical structures with accelerations in opposing directions, giving rise to reversing jets.

The zonal momentum balance can be written as

$$u_t - \beta y v = -\frac{1}{\rho} P_x - uu_x - vv_y - ww_z + \frac{1}{\rho} \tau_z, \quad (1)$$

where subscripts *t*, *x*, *y*, and *z* indicate temporal, zonal, meridional, and vertical derivatives; (*u*, *v*, *w*) are, respectively, the zonal, meridional, and vertical flow; β is the meridional gradient in the planetary vorticity; *y* is the meridional displacement from the equator, which is zero for our analysis of equatorial momentum balance; *P<sub>x</sub>* is the zonal pressure gradient; ρ is the ocean density; and τ is the zonal stress. To derive the layer-averaged zonal momentum balance, the depth-dependent balance (1) is vertically averaged from the top of the layer to the bottom of the layer. Doing this, the layer-averaged zonal momentum balance can be written as

$$\langle u_t \rangle_i - \beta y \langle v \rangle_i = -\left\langle \frac{1}{\rho} P_x \right\rangle_i - \langle uu_x \rangle_i - \langle vv_y \rangle_i - \langle ww_z \rangle_i + \frac{\tau_{i-1} - \tau_i}{\rho(h_i - h_{i-1})}, \quad (2)$$

where  $\langle (\ ) \rangle_i$  is the layer *i* vertical average defined by the layer depths  $z = -h_{i-1}$  and  $z = -h_i$ :

$$\langle (\ ) \rangle_i = \frac{1}{h_i - h_{i-1}} \int_{-h_i}^{-h_{i-1}} (\ ) dz,$$

τ<sub>*i*</sub> is the zonal stress at  $z = -h_i$ , and τ<sub>0</sub> is the zonal wind stress at the ocean surface. In both the depth-dependent (1) and layer-averaged (2) zonal momentum balance, the first term on the lhs is the local acceleration (also referred to as the tendency rate), the second term is zero on the equator, the first term on the rhs is the zonal pressure gradient force, the second, third, and fourth terms are the nonlinear advection terms, and the final term on the rhs is the vertical divergence of the turbulent stress (the horizontal stress divergence is assumed to be negligible). The layers *h<sub>i</sub>* are arbitrary; however, as discussed above, we define *h<sub>1</sub>* as the 28°C isotherm depth (or 25 m, whichever is deeper), *h<sub>2</sub>* as

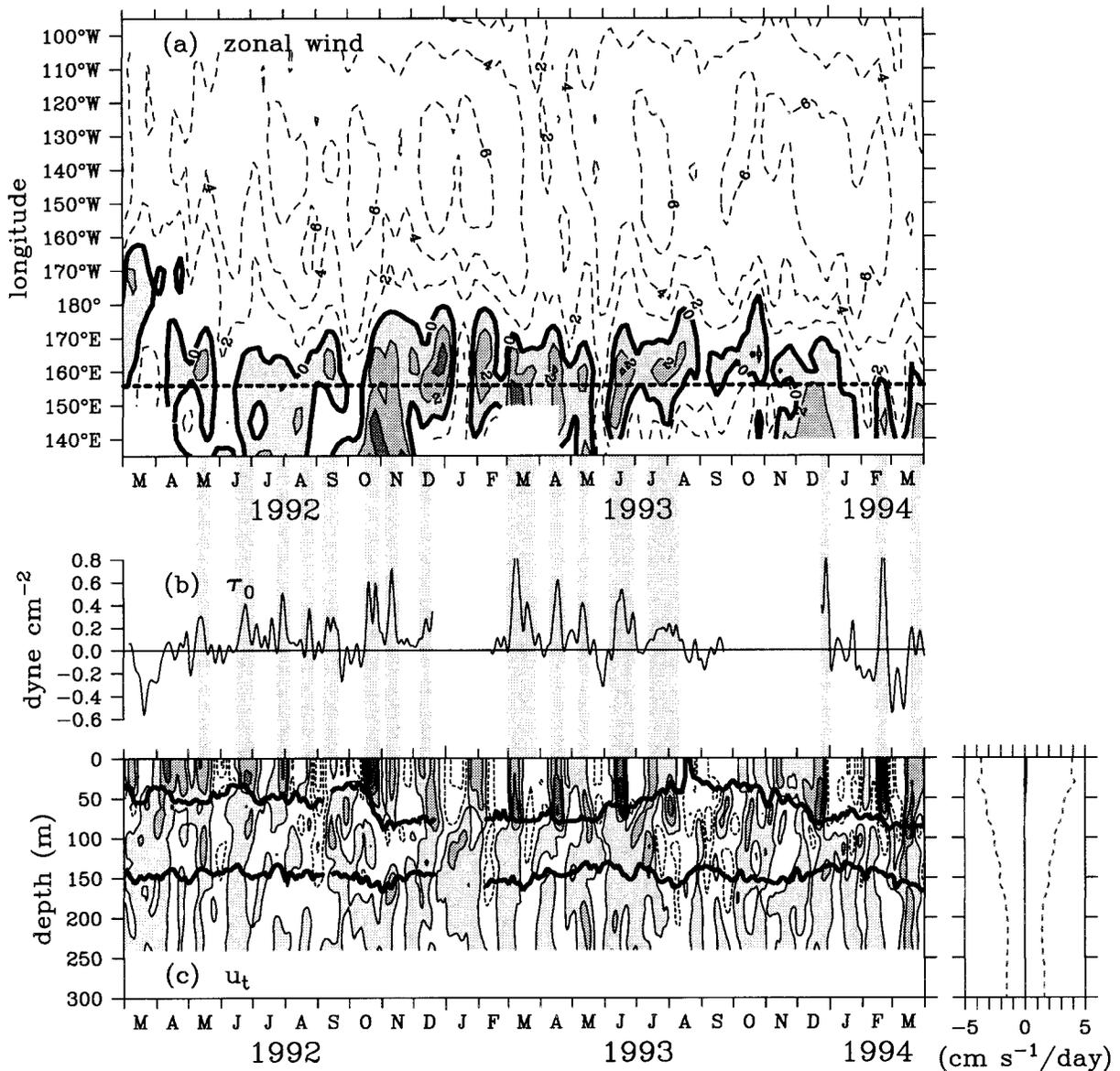


FIG. 5. (a) Time-longitude plot of the 5-day gridded zonal winds along the equator as measured by the TAO array: CI is  $2 \text{ m s}^{-1}$ ; positive values are shaded. (b) Five-day filtered zonal wind stress at  $0^\circ$ ,  $156^\circ\text{E}$ . Large positive (westerly) wind stress events are striped gray. (c) Fifteen-day filtered local zonal acceleration at  $0^\circ$ ,  $156^\circ\text{E}$ : CI is  $2 \text{ cm s}^{-1}/\text{day}$ ; eastward accelerations are shaded. Right panel shows the record length mean and standard deviation envelope of the local acceleration.

the  $20^\circ\text{C}$  isotherm depth,  $h_3$  as 240 m (near the bottom ADCP bin depth), and  $h_0$  as the air-sea interface (0 m).

The pressure gradient force at  $0^\circ$ ,  $156^\circ\text{E}$  shown in Fig. 6a was estimated using the zonal gradient of the 5-day triangular filtered 500-db referenced dynamic height. Although deep flow is not completely motionless (e.g., Firing et al. 1998), the pressure gradient at 500 m is expected to be relatively small near the equator. The zonal gradient of the dynamic height was computed as a “centered” difference between  $0^\circ$ ,  $160.5^\circ\text{E}$  and  $0^\circ$ ,  $154^\circ\text{E}$  for the periods 10 May 1992–20 January 1993

and 13 March–18 April 1993, as the “narrow side difference” between  $0^\circ$ ,  $160.5^\circ\text{E}$  and  $0^\circ$ ,  $156^\circ\text{E}$  for the 14-day period from 19 April through 2 May 1993, and as a “wide side difference” between  $0^\circ$ ,  $165^\circ\text{E}$  and  $0^\circ$ ,  $156^\circ\text{E}$  for all other periods. The pressure gradient estimated with a wide side difference has less high-frequency, localized variability. Consequently, the wind burst pressure gradient response is best analyzed during the well-resolved period between May 1992 and April 1993.

To address the dynamical influence of salinity vari-

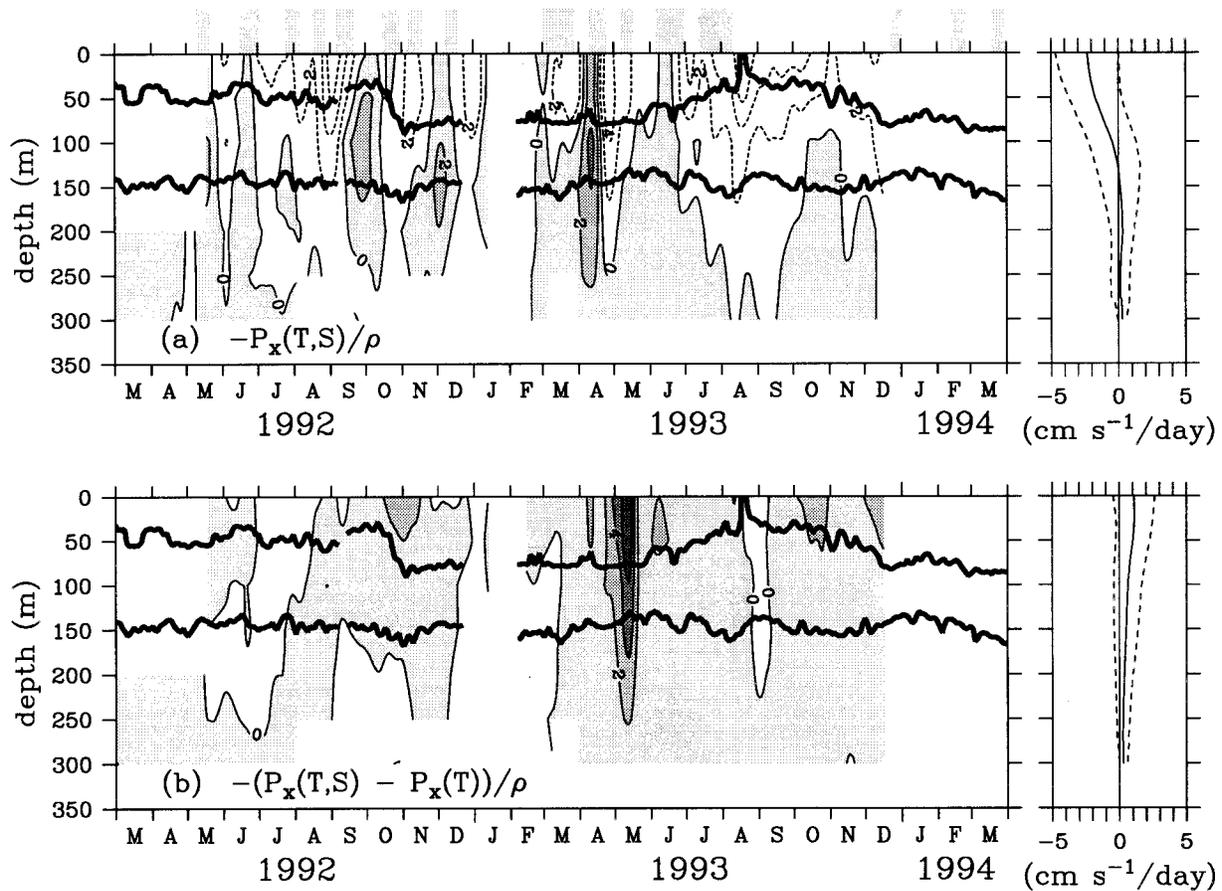


FIG. 6. (a) Fifteen-day triangular filtered zonal force due to the zonal gradient in dynamic height relative to 500 db computed using observed temperature and salinity data. Westerly wind stress events are striped gray as in Fig. 5. (b) The difference between the zonal pressure gradient force computed using temperature and salinity, and computed with temperature only. In both, the CI is  $2 \text{ cm s}^{-1}$  per day and eastward forcing is shaded. The upper and lower thick lines in the left panels are, respectively, the  $28^\circ\text{C}$  and  $20^\circ\text{C}$  isotherms at  $0^\circ, 156^\circ\text{E}$ . Right panels show the mean profile of the time series and its standard deviation envelope.

ability, the pressure gradient force was also computed using only temperature data (salinity was computed from temperature using the Levitus climatological  $T$ - $S$  curves). The difference between the pressure gradient force computed with observed temperature and salinity and the pressure gradient computed using temperature only is shown in Fig. 6b and can be interpreted as the pressure gradient force due to anomalous salinity variability. At the surface, the mean, rms, maximum, and minimum differences are respectively 0.8, 1.7, 10.5, and  $-1.7 \text{ cm s}^{-1}/\text{day}$ , which should be compared to the mean, rms, and maximum and minimum of the surface pressure gradient force (respectively,  $-2.6$ , 3.5, 4.0, and  $-9.9 \text{ cm s}^{-1}/\text{day}$ ). Note that a 3 dyn cm change over 9 degrees of longitude corresponds to a  $2.6 \text{ cm s}^{-1}/\text{day}$  acceleration.

As shown in Fig. 6b, the largest differences in the pressure gradient force estimates tend to occur at the surface, presumably due to the passage of salinity fronts (e.g., October 1992, April–May 1993, and October

1993) (Cronin and McPhaden 1998). However, variability in the salinity maximum between depth 100 and 200 m had a relatively larger effect on the weaker subsurface pressure gradient and local tendency rate. In particular, during the boreal spring and fall of 1993, when the pressure gradient was computed using a wide side difference (between  $0^\circ, 165^\circ\text{E}$  and  $0^\circ, 156^\circ\text{E}$ ), salinity variability often controlled the sign of the pressure gradient in the EUC layer. Finally, we note that in the layer balance, the cross correlations with local tendency rate were substantially lower when the pressure gradient force was computed using temperature only (0.19, 0.23, and 0.11 for the surface, intermediate, and EUC layers, respectively, as compared to 0.35, 0.37, and 0.23).

The local tendency rate (Fig. 5c) and advection terms (Fig. 7) were computed using 5-day filtered daily ( $u$ ,  $v$ ,  $w$ ) velocity data. In particular, the local tendency rate was estimated as a 3-day centered finite difference of zonal velocity at  $0^\circ, 156^\circ\text{E}$ . Zonal advection was computed using a centered finite difference of the zonal

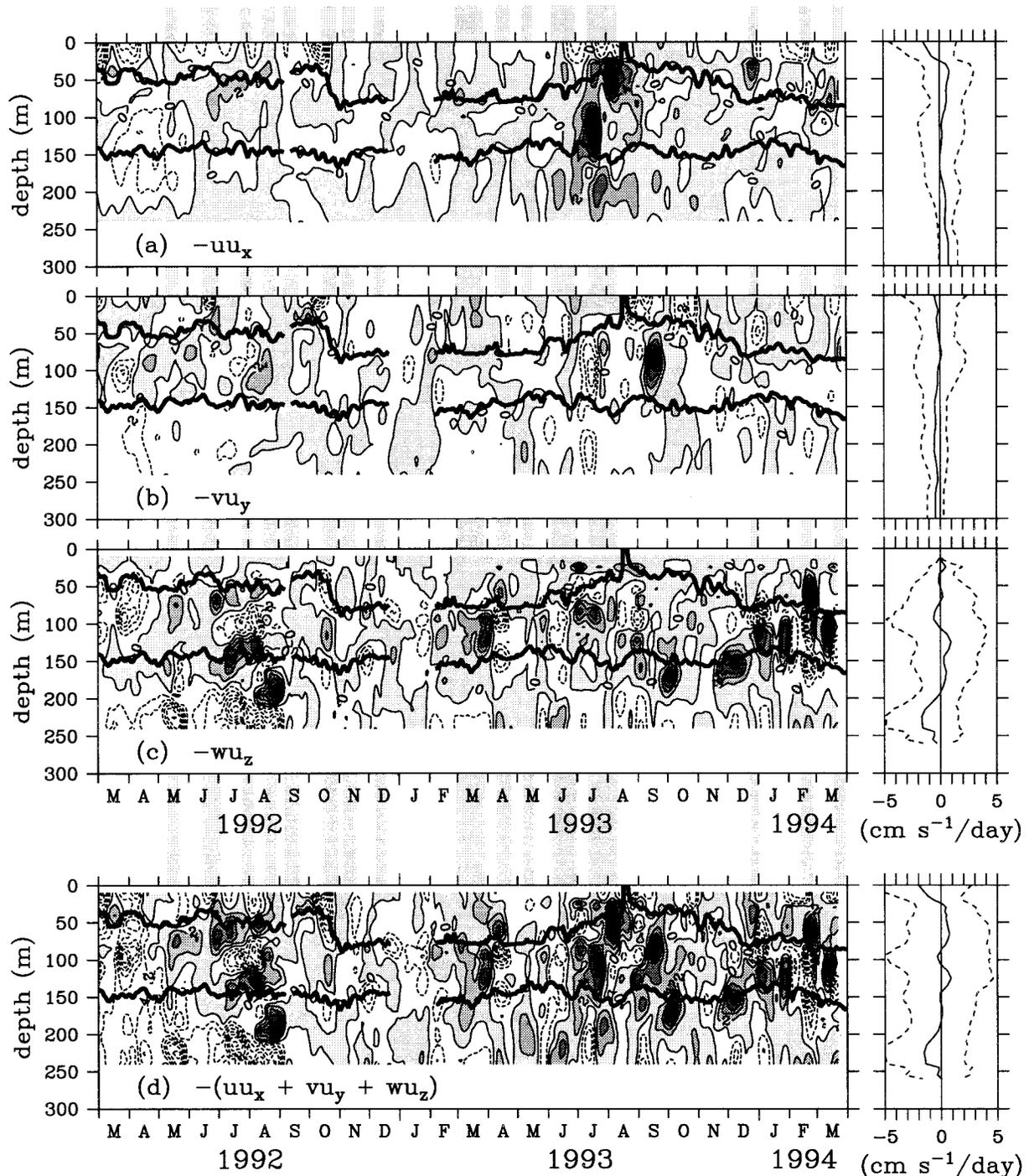


FIG. 7. (a) Zonal, (b) meridional, and (c) vertical advection, and (d) the net 3D advection of zonal flow. The upper and lower thick lines are, respectively, the 28°C and 20°C isotherms; CI is  $2 \text{ cm s}^{-1}/\text{day}$ ; positive advctions (inducing eastward accelerations) are shaded. Westerly wind stress events are striped gray as in Fig. 5. Mean and standard deviation envelopes of the time series are shown in the panels on the right.

velocity ( $0^\circ, 157.5^\circ\text{E}$  minus  $0^\circ, 154^\circ\text{E}$ ) from March 1992 through February 1993 and a side finite difference ( $0^\circ, 157.5^\circ\text{E}$  minus  $0^\circ, 156^\circ\text{E}$ ) from February 1993 through March 1994. Meridional advection was computed using

an upstream finite difference throughout the entire record, and vertical advection was computed using a centered finite difference over the 10-m ADCP bin size.

For the sake of consistency between the layered and

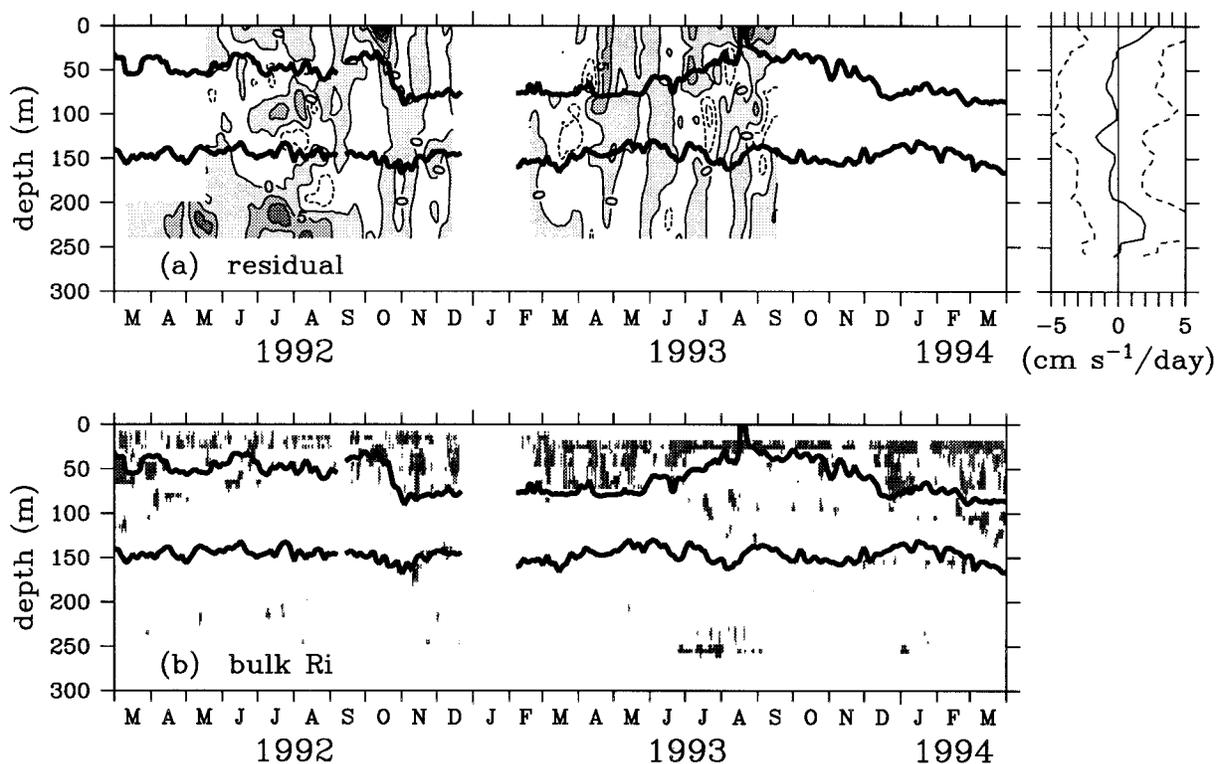


FIG. 8. (a) Residual profile time series: CI is  $5 \text{ cm s}^{-1}/\text{day}$ ; eastward forcing is shaded. The mean profile and standard deviation envelope are shown in the right panel. (b) Bulk Richardson number values less than 1. The upper and lower thick lines are, respectively, the  $28^\circ\text{C}$  and  $20^\circ\text{C}$  isotherms.

depth-dependent analyses, we decompose the zonal stress  $\tau$  into a portion  $\tau_b$  that represents the projection of the wind stress onto the surface layer, and a deviation  $\tau'$  from this wind stress projection:

$$\tau_b = \tau_0 \left( \frac{z + h_1}{h_1} \right) H(z + h_1) \quad \tau' = \tau - \tau_b, \quad (3)$$

where  $H(z + h_1)$  is the Heaviside operator equal to 1 and 0, respectively, above and below  $z = -h_1$ . Thus,  $\tau_b$  is equal to the wind stress  $\tau_0$  at the surface and zero at and below the base of the surface layer, and the vertical divergence of  $\tau_b$  in (1) is equivalent to the wind forcing term in the layer balance (2):  $\tau_{bz}/\rho = \tau_0/\rho h_1$ . This wind stress body force was estimated using hourly wind stress and surface-layer depth time series at  $0^\circ$ ,  $156^\circ\text{E}$  and was subsequently filtered with a 5-day triangular filter and subsampled to once per day.

In general, most terms could be estimated at the  $0^\circ$ ,  $156^\circ\text{E}$  central site. The term  $\tau'_z/\rho$  was estimated as a residual of the zonal momentum balance for periods in which all other terms could be evaluated. The residual, however, also contains the accumulation of sampling and measurement errors and therefore caution is warranted when interpreting it in terms of physics. The depth-dependent residual is shown in Fig. 8 and will be discussed in section 4d.

For presentation purposes, the depth-dependent tendency rate (Fig. 5c), zonal pressure gradient force (Fig. 6), nonlinear terms (Fig. 7), and residual (Fig. 8) were filtered with a 15-day triangular filter (8-day cutoff, 18-day half amplitude). However, in order to analyze processes occurring on wind burst timescales, the layer balance terms shown in Figs. 9–11 were filtered with a 5-day triangular filter. Likewise, all means and standard deviations (Table 2, Figs. 5–8), and cross correlations (Table 3, Figs. 12, 13) were computed using the 5-day triangular filtered time series.

#### 4. Results

Over the record length, the mean tendency rate is near zero both at the surface and at depth (Fig. 5c, Table 2). However, instantaneously, the flow was far from steady, particularly at the surface. Surface currents were strong and variable with frequent changes in direction (Fig. 3), and correspondingly, surface current accelerations were large and variable, with frequent changes in direction (Fig. 5c). Much of this variability can be attributed to direct wind forcing. In contrast to the steady trades found in the central equatorial Pacific, the western equatorial Pacific was characterized by weak trades punctuated by westerly wind

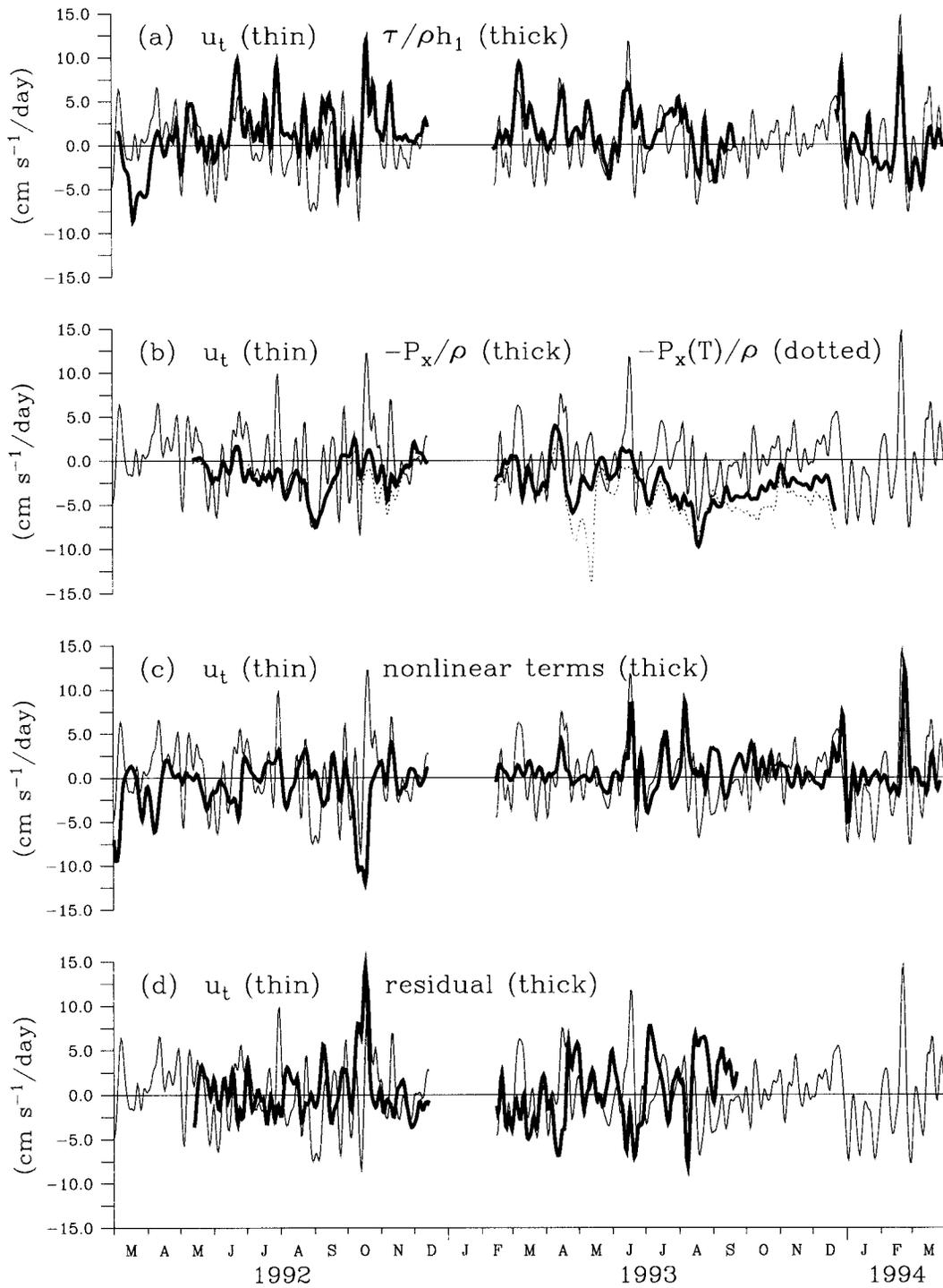


FIG. 9. Surface-layer zonal momentum balance in units centimeters per second per day. The surface-layer tendency rate is repeated as a thin line in each panel, with the following terms shown as a thick line: (a) surface wind forcing; (b) pressure gradient force (the pressure gradient force computed using temperature only is shown as a dotted line); (c) the net zonal, meridional, and vertical advection; and (d) the residual of the surface-layer balance. The surface layer is defined as the layer between the surface interface ( $z = 0$ ) and the depth of the  $28^{\circ}\text{C}$  isotherm (or 25 m, whichever is deeper). All terms are filtered with a 5-day triangular filter.

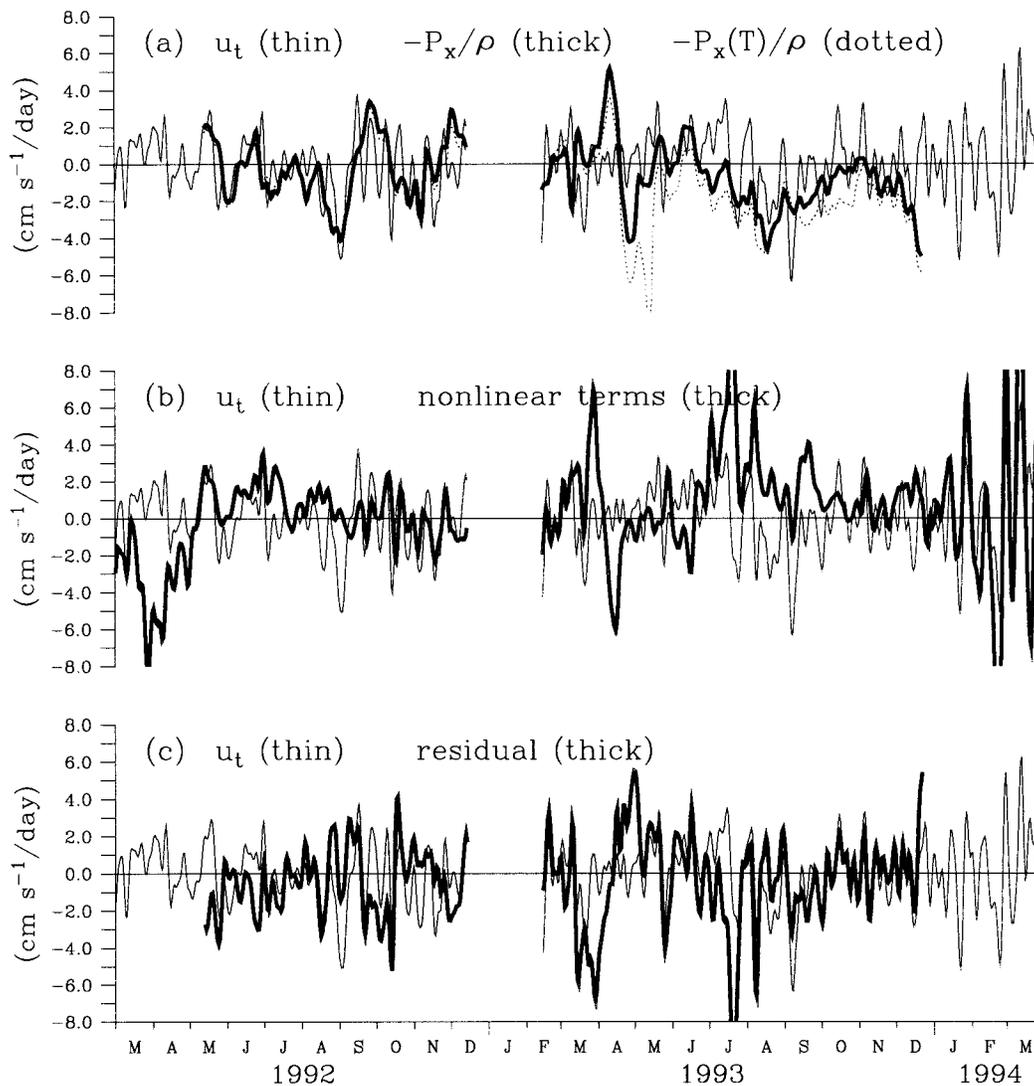


FIG. 10. Intermediate-layer zonal momentum balance in units centimeters per second per day. The intermediate-layer tendency rate is repeated as a thin line in each panel, with the following terms shown as a thick line: (a) pressure gradient force (the pressure gradient force computed using temperature only is shown as a dotted line); (b) the net zonal, meridional, and vertical advection; and (c) the residual of the intermediate-layer balance. The intermediate layer is defined as the layer below 25 m between the 28°C and 20°C isotherm surfaces. All terms are filtered with a 5-day triangular filter.

bursts (Fig. 5a). At 0°, 156°E the record-length mean zonal wind stress was eastward ( $0.08 \text{ dyn cm}^{-2}$ ), and a factor of 3 smaller than the standard deviation of the 5-day filtered zonal wind stress ( $0.25 \text{ dyn cm}^{-2}$ ). Consequently, local wind forcing at 0°, 156°E had relatively frequent changes in direction (Fig. 5b). As shown in Fig. 12 and Tables 2 and 3, this local wind forcing was both highly correlated with and of similar magnitude to the local tendency rate. In fact, of all the terms in the zonal momentum balance, the surface-layer tendency rate was most highly correlated (0.59) with the local wind stress forcing. As a result, surface current anomalies tended to form in the direction of the wind stress approximately 3 days after the peak

wind stress (Fig. 13). Approximately 14 days after the peak wind stress, however, subsurface zonal current anomalies tended to form in a direction opposing the local zonal wind stress. As we will show, these subsurface current anomalies were due in part to the effects of wind forcing on the zonal pressure gradient.

In general, all terms in the momentum balance had large variability and were significantly correlated with the local tendency rate (Fig. 12, Table 3). In particular, the pressure gradient force appears to be important in all three layers, the residual had high correlations with the local tendency rate in the intermediate layer, and nonlinear terms were large and appear to gain importance with depth. We discuss each term in detail below.

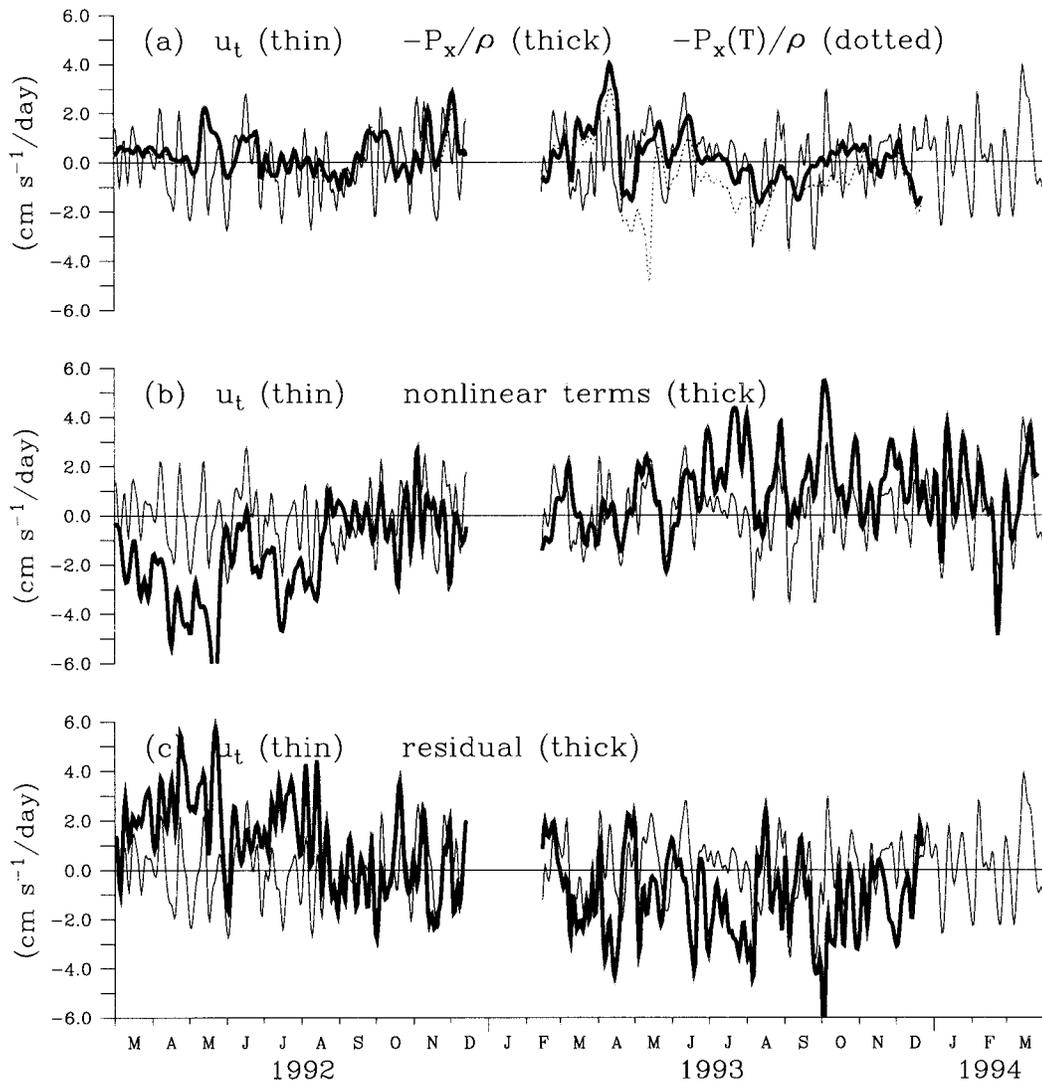


FIG. 11. Same as Fig. 10 but for the EUC layer. The EUC layer is defined as the layer between 20°C and 240 m.

TABLE 2. Means, 90% uncertainty of the means, and standard deviations of 5-day-filtered 0°, 156°E layer-averaged momentum balance. The units of all terms are centimeters per second per day. Values in parentheses indicate mean estimates that are insignificant at the 90% confidence level. The surface layer is defined as the layer between the surface interface ( $z = 0$ ) and the depth of the 28°C isotherm (or 25 m, whichever is deeper). The intermediate layer is the layer below 25 m between the 28°C and the 20°C isotherms, and the Equatorial Undercurrent (EUC) layer is the layer between the 20°C isotherm and 240 m.

	Surface layer		Intermediate layer		EUC layer	
	Mean	Std dev	Mean	Std dev	Mean	Std dev
$u_t$	(0.19 ± 0.42)	3.47	(-0.07 ± 0.27)	1.76	(0.01 ± 0.16)	1.27
$\tau_0/(\rho h_1)$	1.05 ± 0.66	3.09	N/A	N/A	N/A	N/A
$-P_x/\rho$	-2.23 ± 0.77	2.26	-0.57 ± 0.45	1.76	(0.26 ± 0.28)	0.93
$-uu_x$	(-0.23 ± 0.31)	1.84	(0.37 ± 0.53)	1.47	0.36 ± 0.35	1.05
$-vu_y$	0.13 ± 0.13	1.03	(-0.12 ± 0.13)	0.91	(-0.01 ± 0.07)	0.53
$-wu_z$	0.13 ± 0.13	1.07	(0.10 ± 0.29)	2.36	-0.45 ± 0.36	1.35
$-\mathbf{U} \cdot \nabla u$	(0.02 ± 0.33)	2.67	(0.35 ± 0.71)	2.94	(-0.10 ± 0.50)	2.07
Residual	(0.33 ± 0.78)	3.30	-0.45 ± 0.28	2.29	(-0.04 ± 0.50)	2.10

TABLE 3. Cross-correlations between the 5-day filtered layer-averaged local tendency rate and forcing terms. Uncertainties represent the correlation for the null hypothesis at the 90% confidence level. Values in parentheses indicate cross-correlations that are insignificant at the 90% confidence level.

$R_{xy}$ with $u_t$	Surface layer	Intermediate layer	EUC layer
$\tau_0/(\rho h_t)$	$0.59 \pm 0.14$	N/A	N/A
$-P_x/\rho$	$0.35 \pm 0.15$	$0.37 \pm 0.15$	$0.23 \pm 0.13$
$-uu_x$	$(0.03 \pm 0.13)$	$(0.00 \pm 0.13)$	$0.22 \pm 0.11$
$-vu_y$	$0.23 \pm 0.12$	$(0.01 \pm 0.13)$	$0.25 \pm 0.12$
$-wu_z$	$0.23 \pm 0.13$	$0.27 \pm 0.13$	$0.21 \pm 0.11$
$-\mathbf{U} \cdot \nabla \mathbf{u}$	$0.20 \pm 0.13$	$0.22 \pm 0.13$	$0.31 \pm 0.10$
Residual	$(-0.03 \pm 0.17)$	$0.36 \pm 0.14$	$0.22 \pm 0.11$

### a. Wind stress forcing

The surface-layer zonal flow was to a large extent forced by the local winds (Figs. 5, 9, and 12, Tables 2 and 3). However, careful inspection of the surface momentum balance time series (Fig. 9) shows that direct wind-forced Yoshida jet dynamics (accelerations due exclusively to wind stress forcing) occurred primarily during the onset of wind bursts, before the wind stress had reached peak strength. This is also borne out in a lagged correlation analysis of the zonal wind stress and current accelerations (Fig. 13). Peak surface current accelerations occurred approximately 1 day prior to the peak wind stress.

### b. Zonal pressure gradient force

The pressure gradient force was significantly correlated with the local tendency rate in all three layers (Figs. 9–12, Table 3). In particular, during the well-resolved period, there was a near balance between the pressure gradient force and local acceleration in the intermediate layer. Over the record length, the mean surface-layer pressure gradient force was large and westward, opposing the mean eastward wind stress. Indeed, wind stress and lagged pressure gradient force correlations computed over the well-resolved period (Fig. 13c) indicate that approximately 8–10 days after the peak wind stress, a pressure gradient force opposing the wind stress tended to form in the surface and intermediate layers. These pressure gradients generated by WWBs were the primary cause of westward accelerations in the intermediate layer (Fig. 10).

While there were many periods in which the magnitude of the dynamic height gradient increased monotonically from 500 m to the surface, there were also many periods in which the pressure gradient exhibits substantial vertical structure with middepth maxima (Fig. 6a). Note that a nonzero pressure gradient at the reference level (500 db) would bias the pressure gradient at all shallower depths, and therefore would shift the “zero” contour, without changing the vertical structure.

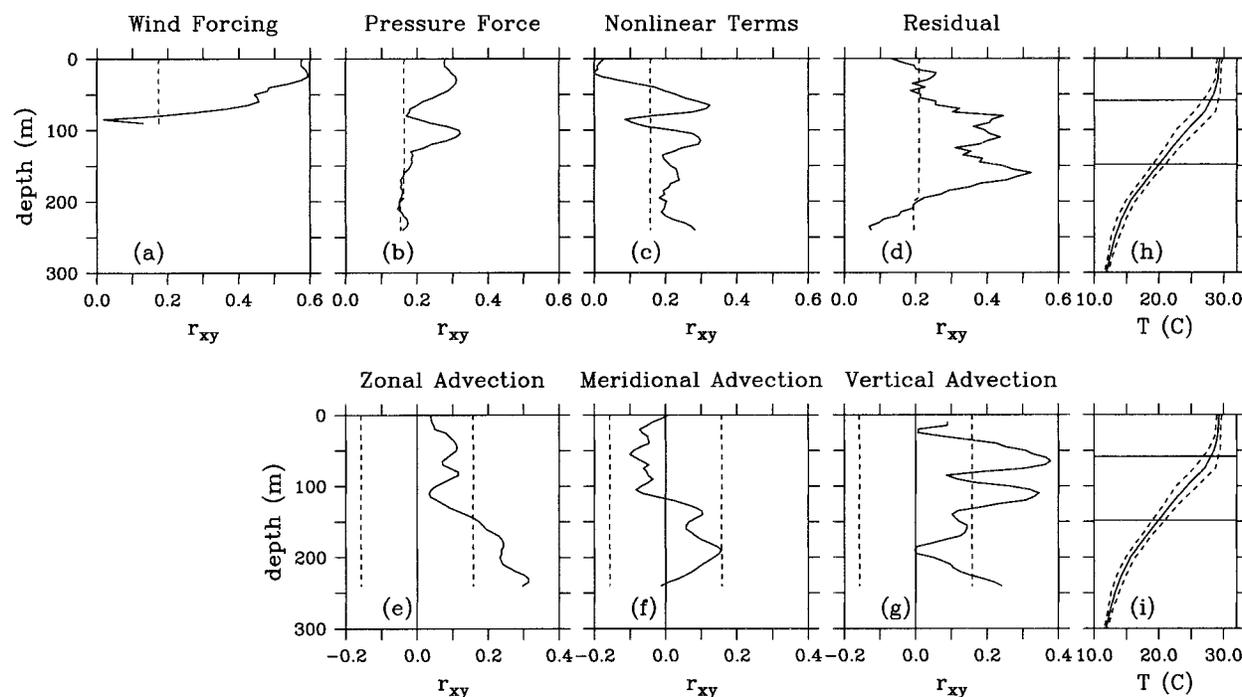


FIG. 12. Cross-correlations between the  $0^\circ$ ,  $156^\circ\text{E}$  depth-dependent local tendency rate and the (a) body force stress divergence (wind forcing), (b) zonal pressure gradient force, (c) nonlinear terms (net zonal, meridional, and vertical advection), (d) residual, (e) zonal advection, (f) meridional advection, and (g) vertical advection. (h) Mean temperature profile and its standard deviation envelope at  $0^\circ$ ,  $156^\circ\text{E}$ . (i) Same as (h). The two horizontal lines in (h) and (i) show the mean depths of the  $28^\circ\text{C}$  and  $20^\circ\text{C}$  isotherms, which define the surface, intermediate, and Equatorial Undercurrent layers. Correlations are computed over the entire record. The dashed lines in (a)–(g) represent the 90% significance levels.

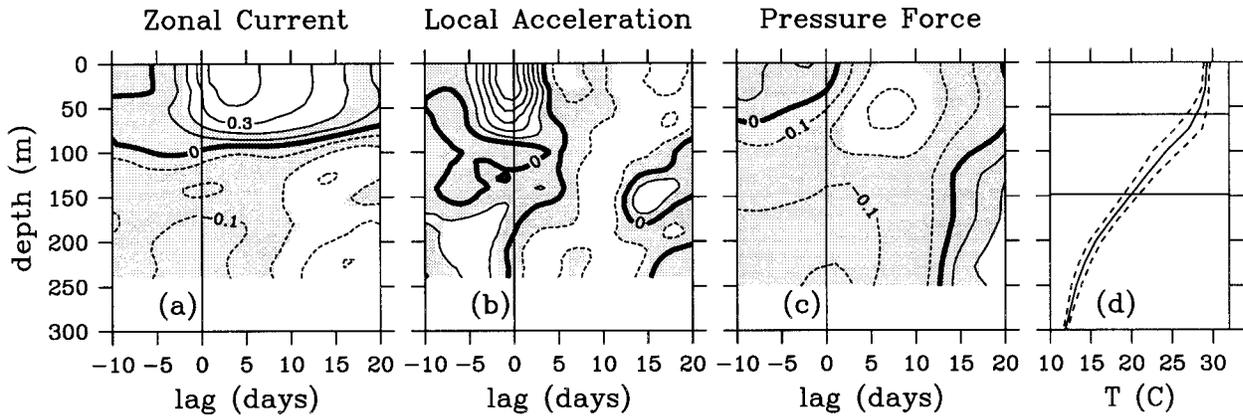


FIG. 13. Lagged correlation between local zonal wind stress and (a) zonal current, (b) local zonal current acceleration, and (c) zonal pressure gradient force as a function of depth at  $0^{\circ}$ ,  $156^{\circ}\text{E}$ . A positive lag implies that wind stress variability occurs prior to the respective variable anomaly. (d) Mean temperature profile and its standard deviation envelope at  $0^{\circ}$ ,  $156^{\circ}\text{E}$ . The two horizontal lines show the mean depths of the  $28^{\circ}\text{C}$  and  $20^{\circ}\text{C}$  isotherms, which define the surface, intermediate, and Equatorial Undercurrent layers. Correlations are computed over the entire record for (a) and (b), and over the period May 1992–April 1993 for (c). Correlations that are not significant at the 90% confidence level are shaded. The CI is 0.1.

### c. Nonlinearities

Nonlinear terms are as much a result of local accelerations having occurred, as they are a cause of the accelerations. Therefore, correlations with the local tendency rate are not required for nonlinearity to be important in the balance. Nevertheless, meridional and vertical advection were significantly correlated with the tendency rate in the surface layer; vertical advection was significantly correlated in the intermediate layer; and zonal, meridional, and vertical advection were all significantly correlated with the local tendency rate in the EUC layer (Table 3). More importantly, zonal and vertical advection were large relative to the local acceleration, particularly in the intermediate and EUC layers (Fig. 7, Table 2).

Vertical shear associated with the reversing jets caused vertical advection to be largest at the depths of the current reversals (Fig. 7c). At the EUC–intermediate layer interface, vertical advection tended to be positive (inducing eastward accelerations) primarily due to upwelling, which brought eastward EUC flow into the intermediate layer. On the other hand, vertical advection at the intermediate–surface layer interface was at times positive and at other times negative due to a variety of vertical velocity and vertical shear combinations. Although some vertical advection events were associated with local Ekman upwelling and downwelling, the strong vertical advection in the EUC layer during April–August 1992 was due to prolonged upwelling that cannot be explained by local Ekman dynamics. In fact, in marked contrast to the correlation analyses shown in Fig. 13, lagged correlations between local zonal wind stress and the nonlinear terms did not show significant correlation patterns and therefore are not shown.

One of the strongest nonlinear events was during July and August 1993, when zonal advection strengthened

the eastward flowing EUC, and vertical advection upwelled the EUC into the intermediate layer. Together with eastward wind forcing in the surface layer, these nonlinearities were largely responsible for the anomalous eastward flow throughout the upper 240 m during the boreal summer of 1993. Similar but less extreme nonlinearities and anomalous eastward flow were also observed during the boreal summer of 1992.

### d. The residual

In principle, the residual (Fig. 8a) should be due to the stress divergence, which deviates from the wind body force. However, since the residual also includes the effects of the pressure gradient at 500 db and the accumulation of computational and sampling errors, it cannot be interpreted unambiguously in terms of physical processes such as vertical mixing. Mixing in the surface layer is due to a combination of wind stirring, nighttime convection, and shear instability, while below the surface layer mixing is due primarily to shear instability. Since vertical mixing within the surface layer may enter both the wind body force term and the residual according to (3), we will focus on subsurface mixing. Shear instability is expected in regions of low Richardson number [ $\text{Ri} = N^2/(u_z^2 + v_z^2)$ ], that is, regions with very high shear ( $u_z^2 + v_z^2$ ) and weak stratification ( $N^2 \sim -g\rho_z/\rho_0$ , where  $g$  is gravity). The bulk Richardson number, shown in Fig. 8b, was estimated using the 5-day filtered current and temperature time series. Note that our estimate of the bulk Richardson number is expected to be larger than an estimate based on high vertical resolution, temporally unfiltered data that would be required to quantify the effects of shear instability on individual mixing events. Qualitatively, however, we found low Richardson numbers predominately in the

surface layer, with occasional low values also in the intermediate and EUC layers. Thus subsurface residuals may well be due in part to vertical mixing.

## 5. Discussion and summary

The western equatorial Pacific during the study period (March 1992 through April 1994) was characterized by weak trade winds punctuated by frequent westerly wind bursts. Consequently, the upper ocean was constantly readjusting to episodic wind stress forcing. Yoshida jet dynamics (accelerations due exclusively to wind stress forcing) were observed in the surface layer above the thermocline during the onset of wind bursts. Consequently, zonal currents lagged zonal wind stresses by about 3 days on average (Fig. 13a). However, typically within 8 days of the peak wind stress, a pressure gradient developed that partially compensated for the wind stress forcing (Fig. 13c). Pressure gradients generated by the wind bursts extended below the surface layer and caused accelerations in the upper thermocline in a direction opposing the wind forcing 8–11 days after the peak wind forcing (Fig. 13). The near balance between the pressure gradient force and the local acceleration in the intermediate layer (Fig. 10) suggests that these pressure gradients were due to equatorial waves. Furthermore, the complex vertical structure in the pressure gradient (Fig. 6a) suggests that high-order baroclinic modes may be important. Quantitative projection of the data onto meridional and vertical mode Kelvin and Rossby waves (e.g., Boulanger and Menkes 1995), however, is beyond the scope of this paper.

The importance of pressure gradients in the readjustment process is consistent with the results from the Richardson et al. (1999) model analysis. They found that baroclinic Rossby and Kelvin waves propagating from the edges, particularly the eastern edge, of the wind patch were primarily responsible for the subsurface westward currents near the eastern edge of the wind patch. Thus, in general, the time it takes for a wind compensating pressure gradient to develop at our study site depends on many factors, including the location of the site relative to the wind patch edges and the propagation speed of the baroclinic Kelvin and Rossby waves generated by the wind burst. For typical WWB fetches of 3000–4000 km and equatorial wave speeds of  $1\text{--}3\text{ m s}^{-1}$ , adjustment within the forcing region would begin within approximately 8–12 days. The 8-day lag between the local wind stress and compensating pressure gradient (Fig. 13) that we found is roughly consistent with this range of expected adjustment times.

Our analysis shows that zonal winds drive slablike flows above the thermocline and indirectly drive counter-currents within the upper portion of the thermocline. Although our layer definitions were based on temperature, salinity stratification can have variability independent of the temperature stratification. In particular, the top of the halocline at times is found to be shallower

than the top of the thermocline. The layer between these two surfaces is referred to as the “barrier layer” (Lukas and Lindstrom 1991; Ando and McPhaden 1997) since it inhibits the mixing of cold thermocline water into the surface layer. Just as the “warm pool” can be defined as the region with surface waters warmer than  $28^{\circ}\text{C}$ , a “fresh pool” can be defined as the region with salinity less than 35 psu. Variability in the location of the fresh pool has been documented by Hénin et al. (1998), Delcroix and Picaut (1998), and Cronin and McPhaden (1998). The “fresh pool” and “warm pool” are not necessarily colocated. Roemmich et al. (1994) show that, if the pressure gradient due to a zonal salinity front within a deep isothermal mixed layer is large enough, it could induce a vertically sheared zonal current within the mixed layer, which could then tilt the zonal salinity front vertically to generate barrier layers. Our analysis shows that there are significant salinity-induced pressure gradients in the surface layer and at the depth of the salinity maximum in the intermediate and EUC layers (Figs. 6b and 11a). Using 14 years of hydrographic data, Kessler (1999) showed that the 150–200 m salinity maximum at  $0^{\circ}$ ,  $165^{\circ}\text{E}$  waxed and waned as the high salinity tongue was advected westward and equatorward from a ventilation region in the southeast Pacific. Our results indicate that this salinity advection has a dynamical effect on the Equatorial Undercurrent.

Reversing jets are due primarily to the interplay between wind forcing and compensating pressure gradients. Nevertheless, nonlinear terms (zonal, meridional, and vertical advection) are large (Figs. 7 and 9–11) and have significant correlations with the local acceleration. As discussed by Philander and Pacanowski (1981) and McCreary (1985), assuming the vertical velocity was due to meridional Ekman convergences and divergences and that the vertical shear took the sign of the wind forcing, one might expect that vertical advection would cause eastward accelerations regardless of the direction of the wind stress. That is, easterlies would tend to upwell the EUC and westerlies would tend to downwell wind-driven eastward surface flow. However, a simple relationship between wind stress and zonal shear does not appear to hold in the highly variable reversing jet regime found in the western equatorial Pacific. Likewise, vertical velocity appears to have large variability due to remote forcing, particularly in the EUC layer during the period April–August 1992. Consequently, at this location, vertical advection is significantly more complex than suggested by Philander and Pacanowski (1981) and McCreary (1985).

Zonal and vertical advection were very large during the boreal summers of 1992 and 1993, when the currents became anomalously eastward throughout the upper 240 m. The record length is too short to determine whether these anomalies were isolated events, a feature of the annual cycle, or a manifestation of interannual variability. Both events were during the termination of El Niño conditions. We note, however, that at  $0^{\circ}$ ,  $165^{\circ}\text{E}$ ,

a site which had over eight years of moored current meter data, anomalous eastward flow throughout the upper water column is evident in the annual mean June and July (Yu and McPhaden 1999). Yu and McPhaden reproduced much of this seasonal anomalous eastward flow using a forced linear analytical model. If the anomalous eastward flow observed at  $0^\circ$ ,  $156^\circ\text{E}$  was in fact a manifestation of the seasonal variability in the western equatorial Pacific, then the importance of nonlinearities in the seasonal cycle would need to be revisited.

Mixing was not directly measured in this analysis, and therefore its role in the dynamics of the current flow has not been highlighted. Nevertheless, mixing is expected to be important, particularly in the surface layer where wind stirring and nighttime convection occur, and at surface and subsurface depths of low Richardson number.

In summary, because of the unique dynamics found on the equator and the strong westerly wind bursts in the western equatorial Pacific, the upper-ocean currents in the western equatorial Pacific are highly variable, changing direction both with time and depth. Using data from the TOGA COARE enhanced monitoring array of TAO and ADCP moorings, nearly all terms in the upper-ocean zonal momentum balance have been evaluated at the central site ( $0^\circ$ ,  $156^\circ\text{E}$ ) for up to two years. The analysis highlights the complex interplay of local and remote wind forcing, and the importance of pressure gradients and nonlinearities as the ocean adjusts to these variable winds of the western equatorial Pacific.

*Acknowledgments.* We gratefully acknowledge R. Lukas (UH) for providing the  $0^\circ$ ,  $154^\circ\text{E}$  and  $0^\circ$ ,  $160.5^\circ\text{E}$  SEACAT temperature and salinity data, and K. Kutsuwada (Tokai U.) for providing ADCP data at  $0^\circ$ ,  $147^\circ\text{E}$  and  $0^\circ$ ,  $154^\circ\text{E}$ . R. Helber (USF) kindly provided his estimate of the vertical velocity and the gap-filled zonal currents at  $0^\circ$ ,  $156^\circ\text{E}$ . D. McClurg (PMEL) assisted in the computation of the dynamic heights and the data availability plots. The authors thank J. Donovan (USF), R. Cole (USF), P. Freitag (PMEL), and the PMEL TAO Project Office for their hard work collecting and processing data. Helpful comments were provided by E. Lindstrom and R. Richardson. RHW was supported by NSF Grant OCE9525912. MFC and MJM were supported by NOAA's Office of Global Programs. This is PMEL contribution 1912.

## APPENDIX

### Salinity Extrapolation

Conductivity (from which salinity is computed) is not a standard measurement on TAO moorings. Thus for the computation of dynamic height at TAO thermistor moorings, typically salinity is estimated from temperature using a climatological temperature–salinity ( $T$ – $S$ ) curve (e.g., from Levitus et al. 1994). However, during

the COARE enhanced monitoring period, over 390 CTD casts were made within 100 km of the equator in the COARE domain and, in addition, several moorings had SEACAT temperature and conductivity sensors (Fig. 2b). While the CTDs had good vertical resolution, they were sparsely sampled in time. Likewise, while the moored SEACATs had good temporal resolution, they were sparsely sampled in the vertical. As described below, these CTD and moored temperature and salinity measurements were blended together using optimal interpolation to create time-dependent  $T$ – $S$  curves at each site along the equator between  $0^\circ$ ,  $154^\circ\text{E}$  and  $0^\circ$ ,  $165^\circ\text{E}$ . These  $T$ – $S$  curves were then used to estimate salinity at the temperature measurement levels below the bottom SEACAT depth for computation of the dynamic height relative to 500 db.

As a first step, at each site, a slowly varying  $T$ – $S$  curve was estimated from CTD data. Typically, near deployments and recovery periods there were 2–20 CTDs and during the COARE IOP there were as many as 75 CTDs taken within a 3-week period near the  $0^\circ$ ,  $156^\circ\text{E}$  mooring. Thus, all CTDs falling within a  $2^\circ \times 4^\circ \times 3$  week window centered at a mooring site were averaged on sigma surfaces. These approximately 6-monthly CTD  $T$ – $S$  curves were then interpolated in time to obtain a gridded time-dependent CTD  $T$ – $S$  curve for each site.

Then, to blend in the buoy SEACAT data, the SEACAT  $T$ – $S$  deviations from the CTD  $T$ – $S$  curve were computed for each day and were mapped onto a temperature grid from  $5^\circ$  to  $29^\circ\text{C}$  using an optimal interpolation (OI) scheme (Bretherton et al. 1976). The temperature grid had  $0.5^\circ\text{C}$  resolution above  $27^\circ\text{C}$  and  $1^\circ\text{C}$  resolution below  $27^\circ\text{C}$ . Optimal interpolation uses a prescribed correlation function to determine how to weight “neighboring” perturbation measurements. In our application, we used a Gaussian time- and temperature-lagged correlation function  $C[(t_i - t_j), (T_i - T_j)]$  of the form:

$$C[(t_i - t_j), (T_i - T_j)] = \exp\left[-\left(\frac{t_i - t_j}{\lambda_t}\right)^2\right] \cdot \exp\left[-\left(\frac{T_i - T_j}{\lambda_T}\right)^2\right],$$

where  $t_i$  and  $t_j$  are times of neighboring measurements,  $T_i$  and  $T_j$  are temperature values of neighboring measurements, and  $\lambda_T$  and  $\lambda_t$  are the correlation temperature and timescales. At zero lags, the autocorrelations were increased by a factor  $1/R_0$ , where  $R_0$  (set as 0.7) is equivalent to the fraction of true variance measured. The temporal correlation scale was set as 20 days. Because variations from the mean  $T$ – $S$  curve above the thermocline are expected to be somewhat independent of variations in and below the thermocline, the temperature correlation scale  $\lambda_T$  was assumed to be of the form:

$$\lambda_T(T) = (2^\circ\text{C}) \tanh\left(\frac{26^\circ\text{C} - T}{1^\circ\text{C}}\right) + 3^\circ\text{C}.$$

Thus, the temperature correlation scale is approximately 1°C above 27°C, and approximately 5°C below 25°C. Surface salinity measurements essentially have no impact on the deep  $T$ - $S$  variations; there, the  $T$ - $S$  curve reverts to the slowly varying CTD  $T$ - $S$  curve.

Finally, the OI-gridded  $T$ - $S$  variations were added to the slowly varying CTD  $T$ - $S$  curves to generate a time series of the  $T$ - $S$  curve for the full temperature range at each site along the equator. These  $T$ - $S$  curve time series were then used to determine salinity at thermistor depths below the bottom SEACAT sensor.

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